

Investigation of Local Electronic, Dielectric and Optical Properties of Individual Composite Tabular AgX Emulsion Microcrystals by Cryo-EELS and LMTO-ASA Techniques

Vladimir Oleshko*, Marc Amkreutz **, and Harald Overhof**

*University of Virginia, Dept. Materials Science & Engineering, Charlottesville, VA 22904-4745, USA

**Universität Paderborn, Fachbereich Physik, 33098 Paderborn, Germany

AgX (X=Br, I, Cl) microcrystals are well known as quantum detectors. The unique character of optical absorption, considerable complexity in the valence-band structure, and remarkable quantum-size effects make them attractive for solid-state research. Here we report local electronic, dielectric and optical properties of composite AgBr core-Ag(Br,I) shells tabular AgX microcrystals studied by cryo-EELS and the *ab initio* linear muffin-tin orbital method (LMTO) in its atomic sphere approximation (ASA) [1]. Cryo-EELS (T= -193°C) within areas of 100 nm in size of the AgBr core was performed on a ZEISS CEM902 computerized electron microscope at U= 80 kV using an integrated energy filter connected to a Kontron IBAS-2000 image analysis system. Fig. 1a compares ϵ_1 , ϵ_2 , and $\text{Im}(-1/\epsilon)$ computed using Kramers-Kronig relations and self-consistent LMTO-ASA. For $E \leq \hbar\omega_p = [(\hbar\omega_p^f)^2 + E_g^2]^{0.5} = 21.6$ eV, $\text{Im}(-1/\epsilon)$ describes interband transitions at 4, 7-8, 15-17 and the volume plasmon at 22 eV. The LMTO-ASA calculations of dielectric parameters have been performed from first principles taking into account only direct transitions. Assuming all the transition matrix elements to be equal, independent of the k-point and the band number of the initial and final electron states, the following expressions for ϵ_1 and ϵ_2 were derived [2]:

$$\epsilon_1 = 1 + V_1 \sum_{\mathbf{P}} \int_{\text{BZ}} \frac{1}{[E_{j'}(\vec{k}) - E_j(\vec{k})]^2 [E_{j'}(\vec{k}) - E_j(\vec{k}) - \hbar\omega]} d^3k \quad (1)$$

$$\epsilon_2 = \frac{V_2}{\omega^2} \sum_{\mathbf{j}, \mathbf{j}'} \int_{\text{BZ}} \delta(E_{j'}(\vec{k}) - E_j(\vec{k}) - \hbar\omega) d^3k \quad (2)$$

Here fitting parameters $V_1 = \frac{2\hbar^4 e^2}{\pi^2 m^2} |M|^2 = 1/25$ and $V_2 = \frac{2\hbar^2 e^2}{\pi m^2} |M|^2 = 1/95$, respectively, with M being

the constant transition matrix element. The needed energies have been taken from the self-consistent LMTO-ASA energy band calculations (Fig. 1b). Although ϵ_2 and $\text{Im}(-1/\epsilon)$ fall to zero at 4.3 eV (the direct band gap of AgBr) due to limitations of the model, the maxima of ϵ_2 at 9 eV, 11 eV and 13 eV satisfactorily fit to the intensive composite band at 7-14 eV (with maxima at 8 eV and 10 eV) assigned to the proposed exciton transitions at X point. The computed minimum of ϵ_1 at 14 eV fits well to the experimental minimum of ϵ_1 , similar to the minimum of $\text{Re}(1/\epsilon)$ at 13 eV. Moreover, the crossings of $\text{Re}(1/\epsilon)$ and $\text{Im}(-1/\epsilon)$ at 8 eV coincide as well as of $\text{Re}(1/\epsilon)$ and ϵ_2 at 22 eV, where $\text{Im}(-1/\epsilon)$ has its maximum and ϵ_1 has a local minimum. Furthermore, there is the shoulder of ϵ_2 at 15-16 eV superimposed with the previous band and the corresponding maximum of $\text{Im}(-1/\epsilon)$. For $E > 25$ eV, both computed and experimental curves closely agree. Again, the features at 4-5 eV cannot be computed due to having no direct transitions at this energy. Fig. 1c presents an overview plot of the projected density of states (DOS) splitted into s-, p- and d-orbital momentums of the projected Ag and Br atoms. One can see the Br s-state at -15 eV and the Br p-states and Ag d-states mixing between -5-0 eV. Strong mixing of the Ag s-, p- and valence d-states and the Br valence p- and d-states occur above the Fermi

level (Fig. 1c). The first conduction band at 3-7 eV seems to be mainly a mixture of the Ag s-states and the Br p-states, while the second and the third ones are due to complicated mixing of the Br d- and p-states with the Ag p- and d-states. Fig. 1d shows the local complex refractive index, $N = \epsilon^{1/2} = n+ik$, derived from EELS data and from LMTO-ASA calculations. Here the index of refraction, $n = \{0.5[(\epsilon_1^2 + \epsilon_2^2)^{0.5} + \epsilon_1]\}^{0.5}$, and the extinction coefficient, $k = \{0.5[(\epsilon_1^2 + \epsilon_2^2)^{0.5} - \epsilon_1]\}^{0.5}$ are given. The values of n closely approach 1.0 above 50 eV. The local absorption coefficient, $\mu(E) = E/(\hbar c)[2(\epsilon_1^2 + \epsilon_2^2) - 2\epsilon_1]^{0.5}$ is shown in Fig. 1e. Maxima of the real part of the optical joint DOS (OJDOS), $\text{Re}\{J_1(E)\} = \frac{E\epsilon_2}{0.5\pi E_p^2}$, $E_p = 21.4$ eV at about 5 eV, 10-11 eV and at 18 eV (Fig. 1f) satisfactorily fit to the corresponding maxima of the calculated DOS in Fig. 1c.

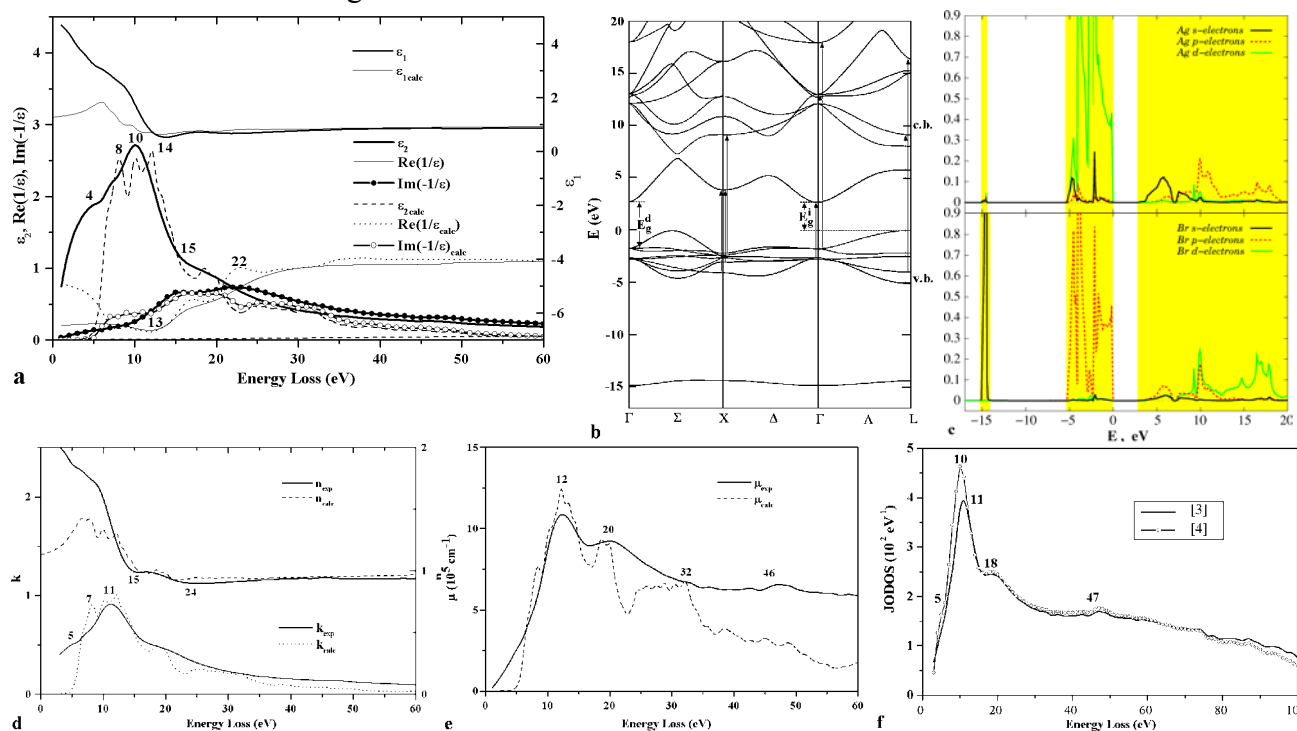


Fig. 1. (a) Dielectric permittivity, $\epsilon = \epsilon_1 + i\epsilon_2$, and energy-loss function, $\text{Im}(-1/\epsilon)$, for an AgBr core of a composite tabular microcrystal vs. LMTO-ASA calculations. (b) Electronic band structure of AgBr calculated by the LMTO-ASA method with a proposed assignment of experimentally observed exciton assisted interband transitions at Γ , L and X points. E_g^d and E_g^i denote energies of the lowest direct and indirect gap, respectively, v.b. is the valence band, c.b. is the conductive band. (c) Projected DOS calculated within the LMTO-ASA approach, splitted into s-, p- and d-orbital momenta of the projected Ag and Br atoms. (d) Refractive index, $N = \epsilon^{1/2} = n+ik$. (e) Absorption coefficient, $\mu(E)$. (f) The real part of OJDOS, $\text{Re}\{J_1(E)\}$, calculated from experimental data using computing programs [3] and [4], respectively.

References

- [1] O.K. Andersen Phys. Rev. B12 (1975) 3060. [2] V. Oleshko, M. Amkreutz, H. Overhof, Phys. Rev., B68 (2003) in press. [3] R.F. Egerton, Electron Energy-Loss Spectroscopy in the Electron Microscope, 2nd Edn. Plenum Press, NY (1996). [4] A.H. Buxbaum, MS Thesis, Northwestern University, Evanston, IL, 1988.