## MEASUREMENTS OF SOLAR INTERNAL ROTATION OBTAINED WITH THE MT. WILSON 60-FOOT SOLAR TOWER

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ABSTRACT. We have obtained estimates of the solar internal rotational velocity from measurements of the frequency splittings of p-mode oscillations. Specifically, we have analyzed a 10-day time series of full-disk Dopplergrams obtained during July and August 1984 at the 60-Foot Tower Telescope of the Mt. Wilson Observatory. The Dopplergrams were obtained with a Na magneto-optical filter and a 244 x 248-pixel CID camera. From the time series we computed power spectra for all of the prograde and retrograde sectoral p-modes from £ = 0 to 200 and for all of the tessaral harmonics up to  $\ell$  = 89. We then applied a crosscorrelation analysis to the resulting sectoral power spectra to obtain estimates of the frequency splittings. From  $\ell = 4$  to  $\ell = 30$  we obtained a mean value of the frequency spitting of roughly 450 nHz (sidereal) in close agreement with most previously published results, while from  $\ell = 40$  to  $\ell = 140$  we obtained a mean value of about 470 nHz. believe that the latter value is slightly higher than the surface rotational splitting of 461 nHz because of possible confusion due to the 42 E. J. RHODES JR. ET AL.

temporal sidelobes introduced by the day/night observing cycle. Confirmation of this possibility will have to await our computation of tesseral power spectra for degrees greater than our current limit of 89. Finally, for degrees between 140 and 200, the frequency splittings are indistinguishable from the surface rotation rate.

## 1. Introduction and Discussion

Several recent studies have obtained estimates of the solar internal equatorial rotation rate from measurements of the frequency splittings of oscillations  $^{1-4}$ . Each of these studies obtained estimates of the rotation rate which were either near or slightly below the surface rotation rate of magnetic features  $^5$ . During the summer of 1984 we obtained solar full-disk Dopplergrams every 40 seconds for up to 11 hours per day on 90 different days. These Dopplergrams were obtained with a magneto-optical filter feeding a CID camera. We have analyzed a 9.5-day subset of these images in order to obtain our own estimates of the internal equatorial rotation rate. Sperical harmonic filtering yielded time series of sectoral harmonics which were Fourier transformed to yield the power spectrum of Figure 1. Figure 2 shows that this spectrum clearly contained frequency shifts induced by solar rotation. Figures 3 and 4 contain some of our results and those of Duvall and Harvey and Libbrecht for comparison. Figure 5 shows all of our results from  $\ell$  = 4 to  $\ell$  =200.

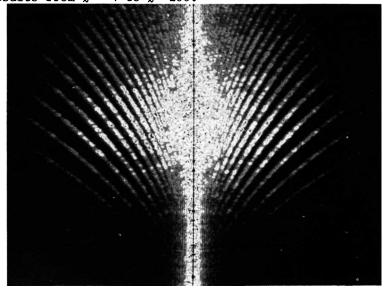


Figure 1. Two-dimensional,  $\ell$ - $\nu$ , power spectrum computed from 9.5-day time series of Sectoral harmonic amplitudes. The vertical black line in the center is the temporal frequency axis, while the x-axis ranges from  $\ell$  = -200 at the left to  $\ell$  = 200 at the right. The sharp ridges of observed power are due to the solar p-modes. At least 23 ridges can be distinquished in each half of the spectrum.

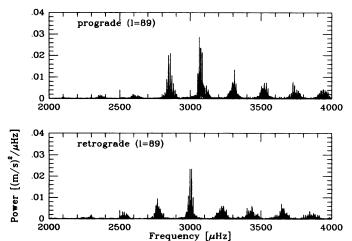


Figure 2. Two vertical slices through the power spectrum shown in Figure 1. Here the prograde and retrograde sectoral harmonic peaks corresponding to a degree of 89 are displayed as a function of the absolute temporal frequency,  $\nu$ , in mHz. The prograde peaks can be clearly seen to have higher frequencies than do the corresponding retrograde peaks in the lower panel. Since the total prograde-minus retrograde frequency splitting is equal to 2 x 89 x  $\Delta\nu$ , where  $\Delta\nu$  is the rotational frequency, it is clear that  $\Delta\nu$  can be computed easily by dividing the measured splitting by 188. The sidelobes present in each peak are due to both the temporal window function and to spatial leakage from adjacent degrees.

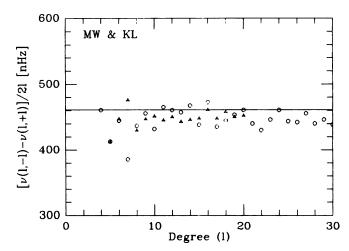


Figure 3. Comparison of our rotational frequency splittings obtained from a cross-correlation of the prograde and retrograde portions of Figure 1 (shown as the circles) with the corresponding measurements of Libbrecht<sup>4</sup> (shown as the triangles). Both sets of splittings appear to lie systematically below the surface rotational rate of magnetic features of 461 mHz (sidereal)<sup>5</sup> which is drawn as the solid line.

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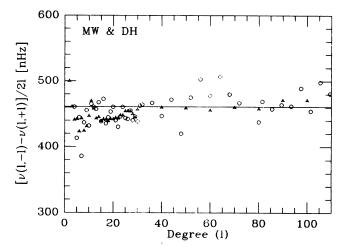


Figure 4. Comparison of our rotational frequency splittings (the circles) with those of Duvall and Harvey (the triangles) for  $\ell=4$  to  $\ell=100$ . Below  $\ell=40$  both sets of results lie systematically below the surface rotation rate (again shown as the solid line). The relatively large scatter of ourvalues above  $\ell=44$  may be due to the shorter duration of our dataset (9.5 days versus 17 days for Duvall and Harvey) or it may be due to increased contamination from the spatial sidelobes from adjacent degrees.

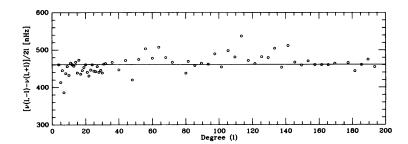


Figure 5. Our rotational splittings extending from  $\ell$  =4 to  $\ell$  =200. As shown in Figures 3 and 4, our values lie below the surface rotation rate of magnetic features for  $\ell$  = 4 to 42, while they average 470 nHz between  $\ell$  = 42 and  $\ell$  = 140. Above  $\ell$  = 140 our results are indistinquishable from the surface rate of 461 nHz.

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