

# PRECESSION AND STELLAR KINEMATIC PARAMETERS FROM THE PROPER MOTIONS OF THE AGK3U

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**ABSTRACT.** Precessional corrections and stellar kinematic parameters have been derived from the proper motions of 142,171 stars in the updated version of AGK3 (AGK3U). The sky has been divided into 1332 small areas, in which mean proper motions for each component are formed for different magnitude intervals, with and without exclusion of fast stars. Solutions were performed for various kinds of the mean proper motions. Maximum-likelihood algorithm is used to take into account the ellipsoidal distribution of residual velocities. The results for precessional correction are,  $\Delta n = 0''.47 \pm 0''.03$  /cy,  $\Delta k = -0''.35 \pm 0''.02$  /cy, and the results for Oort's constants are,  $A = 9.5 \pm 1.5$  km/s/kpc,  $B = -7.5 \pm 1.5$  km/s/kpc.

## 1. INTRODUCTION

From the proper motions of the preliminary version of AGK3, Dieckvoss (1967) has obtained estimates for precessional corrections  $\Delta n = 0''.51 \pm 0''.01$  /cy,  $\Delta k = -0''.32 \pm 0''.01$  /cy, and galactic rotation constants  $A = 14.8 \pm 0.6$  km/s/kpc,  $B = -11.3 \pm 0.6$  km/s/kpc. Asteriadis (1977) has determined the four quantities  $\Delta n = 0''.44 \pm 0''.02$  /cy,  $\Delta k = -0''.36 \pm 0''.02$  /cy,  $A = 16.1 \pm 1.9$  km/s/kpc,  $B = -9.0 \pm 1.9$  km/s/kpc, from the published AGK3 proper motions.

In this investigation, we aimed to derive the precessional corrections and the constants of galactic rotation on the basis of the AGK3U proper motions. In addition to the improvement of the proper motions, in the present work, we employed a new reduction method, namely maximum-likelihood algorithm, which takes into account both the ellipsoidal distribution of residual velocities and observational errors (Wei, 1987).

## 2. THE MATERIAL

AGK3U is an updated version of AGK3, which improved the AGK3 positions and proper motions using the observations of the Palomar 'Quick V' survey made for the construction of Hubble Space Telescope Guide Star Catalogue (Bucciarelli et al, 1992). It provides new positions and proper motions for 170,464 stars north of  $-2.5$  degree declination to the limiting magnitude 13.3 on magnetic tape. The positions have a mean error of  $0''.167$  at an

average epoch of 1950.62 and the proper motions have a two dimensional mean error of  $0''.82$  /cy.

For each star the magnetic tape contains, besides the usual information (position, proper motion, magnitude, spectral type etc.), an error flag to indicate if linear model for determination of proper motion failed. Of the 170,464 AGK3U stars, we rejected 15,745 stars with an error flag, 5,735 OB stars and 6,813 stars brighter than  $8^m$ .

### 3. DIVISION OF THE SKY AND FORMATION OF NORMAL PROPER MOTIONS

The sky has been divided into 1332 areas of  $4^\circ \times 4^\circ$  and normal proper motions formed within the areas, so that the computations can be made more economically. The loss of accuracy in the final results caused by such a division is negligible, according to Asteriadis(1977).

There are many stars with large proper motions in AGK3U. More than one tenth of the stars have a total proper motion larger than  $6''$  per century. The percentage of large motions decrease with magnitude. ( 16 percent in magnitude interval  $8.0 \leq m_{pg} < 9.0$ , while 6 percent at  $m_{pg} \geq 11.0$ . In order to examine the effect of large proper motions in the solutions, we have formed the normal proper motion in the following ways:

- 1,  $\mu < 6''/\text{cy}$ , where  $\mu = ((\mu_\alpha \cos \delta)^2 + \mu_\delta^2)^{1/2}$
- 2,  $\mu < 12''/\text{cy}$
- 3, all the stars.

Solutions have been performed for the cases in different magnitude intervals and the whole material.

### 4. NUMERICAL METHOD

In the classical analyses of proper motions by the least square method, one has to assume that the combination of observational error and peculiar velocity can be treated as a single variable. In the straightforward case of analyses of a single component, the variance to be assigned to a star should depend on its position due to the assumed ellipsoidal distribution of peculiar velocities even if the observational error is the same for all of the stars. In the case of a combined solution of proper motions of two coordinates, not only will the peculiar velocity components depend on the position, but they will also be correlated(Murray, 1983). In order to overcome this difficulty, a maximum-likelihood algorithm has been developed(Wei, 1987).

It is found empirically that the axes of the velocity ellipsoid are almost the same as the Galactic-coordinate axes(Mihalas and Binney, 1982), so it is reasonable to assume that the Galactic components of the residual  $V$  ( $v_1, v_2, v_3$ ) for an individual stars are sample from Gaussian distributions with zero means and variance  $\sigma_i^2$ , and they are independent of each other.

Let  $\Delta(\Delta_1, \Delta_2)$  be the difference between the observed proper motions and expected proper motions, then such residuals can be expressed in the form of a linear function of  $v_i$ :

$$\Delta_j = \sum_{i=1}^3 a_{ij} V_i$$

It is easy to show that  $\Delta_1, \Delta_2$  must be sampled from a general bivariate Gaussian distribution, and hence for all stars the likelihood function  $L$  is of all of the peculiar velocity components obtained in a practical analysis is given by the product:

$$L = \prod (2\pi)^{-1} |M|^{-1/2} \exp(-\frac{1}{2} \Delta M^{-1} \Delta^T)$$

Taking the observational variance  $\sigma_{\epsilon_j}^2$  into account, we can write for the covariance tensor of the residuals:

$$M = \begin{pmatrix} \sigma_{\epsilon_1}^2 + \sum_{i=1}^3 a_{i1}^2 \sigma_i^2 & \sum_{i=1}^3 a_{i1} a_{i2} \sigma_i^2 \\ \sum_{i=1}^3 a_{i1} a_{i2} \sigma_i^2 & \sigma_{\epsilon_2}^2 + \sum_{i=1}^3 a_{i2}^2 \sigma_i^2 \end{pmatrix}$$

The residuals of each star depend on the nature of the adopted model for the kinematic properties of the galaxy. Following the conventional practice, we have:

$$\begin{aligned} \Delta_1 = \mu_\alpha \cos \delta & - f(X \sin \alpha - Y \cos \alpha) - \omega_1 \cos \alpha \sin \delta - \omega_2 \sin \alpha \cos \delta + \omega_3 \cos \delta \\ & + P(\cos 2l \cos b \cos \phi + \frac{1}{2} \sin 2l \cos 2b \sin \phi) \\ \Delta_2 = \mu_\delta & - f(X \cos \alpha \sin \delta + Y \sin \alpha \sin \delta - Z \cos \delta) + \omega_1 \sin \alpha - \omega_2 \cos \alpha \\ & + P(\cos 2l \cos b \sin \phi - \frac{1}{2} \sin 2l \sin b \cos \phi) \end{aligned}$$

where  $(\mu_\alpha \cos \delta, \mu_\delta)$  are observed proper motions. So the log likelihood function:

$$l = -\frac{1}{2} \sum (\ln |M| + \Delta M^{-1} \Delta^T) + const.$$

is a function of the unknowns  $\sigma_i, X, Y, Z, \omega_i, P, (i=1,3)$ .

Maximum-likelihood estimates of the parameters are those for which the function  $-l$  is a minimum. It could be shown that the covariance matrix of the estimated parameters is the Hessian matrix of  $-l$ :

$$H = - \begin{pmatrix} \frac{\partial^2 l}{\partial x_1^2}, & \frac{\partial^2 l}{\partial x_1 \partial x_2}, & \dots, & \frac{\partial^2 l}{\partial x_1 \partial x_n} \\ \frac{\partial^2 l}{\partial x_1 \partial x_2}, & \frac{\partial^2 l}{\partial x_2^2}, & \dots, & \frac{\partial^2 l}{\partial x_2 \partial x_n} \\ \vdots & \vdots & \dots & \vdots \\ \frac{\partial^2 l}{\partial x_1 \partial x_n}, & \frac{\partial^2 l}{\partial x_2 \partial x_n}, & \dots & \frac{\partial^2 l}{\partial x_n^2} \end{pmatrix}$$

The error of unit weight is:

$$\sigma_0^2 = -2l_0/N$$

where  $N$  is the number of degrees of freedom.

### 5. RESULTS AND DISCUSSION

The results for Solar motion  $(X, Y, Z)$ , the components of angular vector  $(\omega_1, \omega_2, \omega_3)$  and Oort's constant  $P$  are listed in Table 1 to 3, including the precessional corrections  $\Delta n, \Delta k$

Table 1, Results from all usable stars

Mag.	10-11	10-13	11-13	8-10	8-11	8-13	8-9	9-10	9-11	9-13
$\sigma_0$	123	123	176	167	130	136	116	145	141	120
$\sigma_1$	77 2	77 2	137 4	99 3	72 2	67 2	205 6	101 3	75 2	68 2
$\sigma_2$	67 3	67 3	97 5	84 3	67 3	62 2	157 7	93 3	69 3	62 2
$\sigma_3$	51 2	51 2	72 4	50 2	46 2	45 2	78 4	47 2	46 2	44 2
X	0 3	0 3	5 5	7 3	4 3	4 3	11 6	10 3	1 3	2 3
Y	-136 4	-136 4	-102 6	-205 4	-180 3	-171 3	-236 9	-216 5	-168 4	-160 3
Z	124 2	124 2	104 4	158 3	147 2	142 2	183 5	146 3	137 2	134 2
$\omega_1$	16 3	16 3	12 6	16 4	14 3	13 3	11 8	38 4	17 3	17 3
$\omega_2$	-47 3	-47 3	-38 5	-42 3	-45 3	-44 3	-44 6	-41 3	-49 3	-45 3
$\omega_3$	-49 2	-49 2	-50 4	-38 3	-41 2	-42 2	-35 5	-34 3	-42 2	-44 2
$\Delta n$	50 3	50 3	41 5	45 3	48 3	47 3	46 6	49 4	53 3	49 3
$\Delta k$	-41 3	-41 3	-44 5	-30 3	-34 3	-35 2	-29 6	-14 3	-33 3	-35 3
P	19 4	19 4	5 6	17 4	20 3	20 3	2 8	43 4	21 3	20 3
Q	-18 4	-18 4	-14 6	-18 5	-16 4	-15 3	-13 9	-44 5	-20 4	-20 3

Table 2, Results from star with  $\mu < 12''$  /cy

Mag.	10-11	10-13	11-13	8-10	8-11	8-13	8-9	9-10	9-11	9-13
$\sigma_0$	113	113	119	127	124	122	133	124	121	108
$\sigma_1$	54 2	54 2	95 3	60 2	49 2	47 2	114 4	68 2	51 2	49 2
$\sigma_2$	50 2	50 2	78 4	58 2	49 2	46 2	91 4	61 2	50 2	47 2
$\sigma_3$	44 2	44 2	56 3	45 2	42 2	41 2	60 3	46 2	42 2	41 2
X	-2 2	-2 2	2 4	-4 3	-4 2	-3 2	-3 4	-5 3	-4 2	-3 2
Y	-108 3	-108 3	-82 5	-149 3	-134 3	-128 2	-171 5	-140 3	-126 3	-121 3
Z	104 2	104 2	92 3	123 2	117 2	114 2	132 3	119 2	114 2	111 2
$\omega_1$	13 3	13 3	9 4	8 3	8 2	8 2	6 5	10 3	10 2	10 2
$\omega_2$	-44 2	-44 2	-39 4	-42 2	-43 2	-43 2	-40 4	-43 3	-44 2	-44 2
$\omega_3$	-50 2	-50 2	-56 3	-39 2	-42 2	-43 2	-37 3	-39 2	-43 2	-45 2
$\Delta n$	47 2	47 2	41 4	44 2	45 2	45 2	41 4	45 3	46 2	46 2
$\Delta k$	-43 2	-43 2	-51 4	-35 2	-38 2	-39 2	-34 4	-34 3	-38 2	-40 2
P	18 3	18 3	7 4	16 3	18 3	18 2	13 5	16 3	18 3	18 2
Q	-15 3	-15 3	-10 5	-9 3	-9 3	-9 3	-7 5	-12 4	-12 3	-12 3

Table 3, Results from star with  $\mu < 6''$  /cy

Mag.	10-11	10-13	11-13	8-10	8-11	8-13	8-9	9-10	9-11	9-13
$\sigma_0$	115	115	100	103	103	102	105	113	107	102
$\sigma_1$	43 2	43 2	72 3	48 2	39 1	37 1	92 3	51 2	40 2	38 1
$\sigma_2$	40 2	40 2	61 3	44 2	39 1	36 1	64 3	47 2	40 2	37 1
$\sigma_3$	40 1	40 1	49 3	38 1	38 1	38 1	47 2	39 2	39 1	39 1
X	-6 2	-6 2	-1 3	-12 2	-10 2	-10 2	-12 3	-12 2	-10 2	-9 2
Y	-78 2	-78 2	-62 4	-99 2	-91 2	-88 2	-111 4	-95 3	-88 2	-84 2
Z	84 2	84 2	79 2	92 2	90 1	88 1	96 3	90 2	88 1	87 1
$\omega_1$	7 2	7 2	3 3	2 2	3 2	4 2	-1 4	5 2	5 2	5 2
$\omega_2$	-44 2	-44 2	-40 3	-43 2	-44 2	-44 2	-41 3	-44 2	-45 2	-44 2
$\omega_3$	-48 1	-48 1	-55 2	-36 1	-39 1	-41 1	-35 2	-36 2	-41 1	-42 1
$\Delta n$	46 2	46 2	41 3	43 2	45 2	45 2	41 3	45 2	46 2	45 2
$\Delta k$	-44 2	-44 2	-53 3	-35 2	-37 2	-39 2	-36 3	-33 2	-38 2	-39 2
P	17 2	17 2	8 4	15 2	16 2	16 2	11 4	15 3	17 2	17 2
Q	-8 2	-8 2	-3 4	-2 3	-3 2	-5 2	1 4	-6 3	-6 2	-6 2

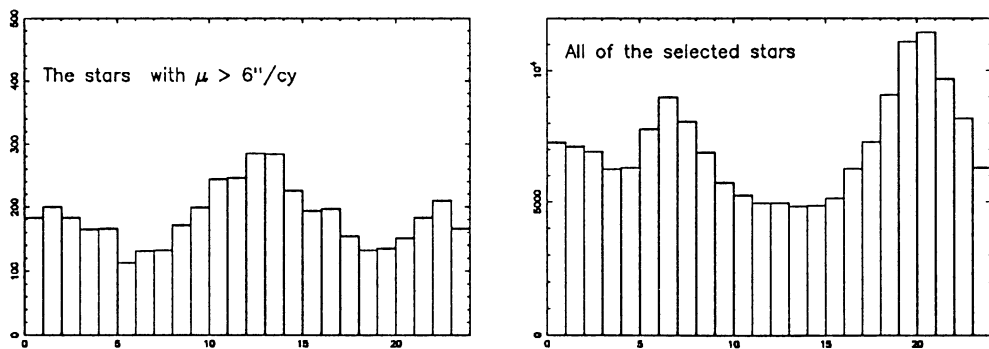


Figure 1, Distribution of AGK3U stars in R. A.

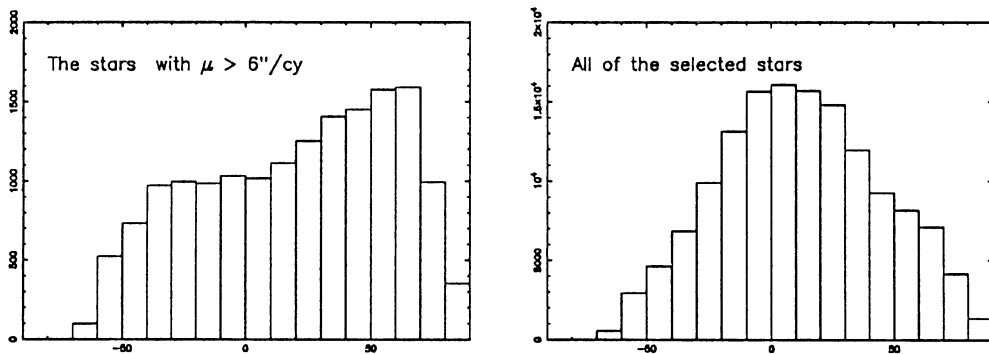


Figure 2, Distribution of AGK3U stars galactic latitude

and Oort's constant  $Q$  derived from  $(\omega_1, \omega_2, \omega_3)$ . The main results can be summarized as follow:

1), *magnitude equation*: No significant magnitude equation was detected. Only the result of solar motion  $(x, y, z)$  varies in different magnitude intervals, and it is resulted from that brighter stars, generally speaking, are nearer to the Sun. But the dependence of solar motion on magnitude interval implies that it might be questionable to use unit parallax factor for all stars.

2), *effect of large proper motions*: There is significant difference in the results of stellar kinematic parameters in the solutions of all stars and with exclusion of large proper motions. This is different from the conclusion of Asteriadis(1977). In order to explain this disagreement, we plotted the  $\alpha$  distribution and  $b$  distribution of all of the selected stars and large proper motion stars (Figure 1 and 2). From these distributions, it can be concluded that the large proper motion are mainly caused by solar motion.

3), *precessional corrections*: The results for precessional corrections are very stable in different solutions. And  $\Delta n = 0''.47 \pm 0''.03$  /cy,  $\Delta k = -0''.35 \pm 0''.02$  /cy resulting from all of the selected stars are in good agreement with the results obtained by Asteriadis(1977) from AGK3.

4), *galactic rotation constants*: The results for Oort's constants are confusing. Oort's constants become quite small after the exclusion of large proper motions. Since the distributions shown in figure 1 and 2 express that the large proper motions contribute very little to the determination of Oort's constants, the dependence of Oort's constants on the fast stars must be resulted from the disadvantage that all of the stars are distributed in north sky.

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