

# A METHOD FOR SOLVING POISSON'S EQUATION FOR SYSTEMS WITH AXIAL SYMMETRY

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In computer experiments on the dynamics of stellar systems special methods are often required for the computation of the forces to keep the problem manageable. For collisionless systems a method based on solving Poisson's equation with Fourier-series expansions has been used with success (Miller *et al.*, 1970; Hohl, 1973). The systems in these studies consist of a large number of particles moving on a rectangular grid.

For three-dimensional systems with axial symmetry a similar method can be used, based on the expansion of density and potential in Legendre polynomials (cf. Prendergast and Tomer, 1970):

$$\varrho(r, \theta) = \sum_{n=0}^N a_n(r) P_n(\cos \theta), \quad (1)$$

$$\psi(r, \theta) = \sum_{n=0}^N b_n(r) P_n(\cos \theta). \quad (2)$$

Poisson's equation,  $\nabla^2 \psi = 4\pi G \varrho$ , with appropriate boundary conditions, supplies the relations between  $a_n$  and  $b_n$  and the potential can be written in the form:

$$\psi(r, \theta) = -4\pi G \sum_{n=0}^N \frac{P_n(\cos \theta)}{2n+1} \left[ r^n \int_r^\infty s^{1-n} a_n(s) ds + \frac{1}{r^{1+n}} \int_0^r s^{2+n} a_n(s) ds \right], \quad (3)$$

where

$$a_n(r) = \frac{2n+1}{2} \int_{-1}^1 \varrho(r, \theta) P_n(\cos \theta) d \cos \theta. \quad (4)$$

(The boundary conditions  $\psi \rightarrow 0$  for  $r \rightarrow \infty$  and  $\psi \rightarrow \text{constant}$  for  $r \rightarrow 0$  have been used). Given a grid in  $r$  and  $\theta$ , the integrals over  $r$  in (3) can be calculated from recurrence relations. The number of operations required to evaluate  $\psi$  from (3) and (4) for a given density distribution is therefore proportional to the number of grid divisions in  $r$ , the number of divisions in  $\theta$  and the number of Legendre polynomials.

We have used this scheme to study the evolution of a system of 4000 particles (similar to the systems studied by Gott, 1973); its numerical behavior is quite satisfactory.

## References

Gott, J. R.: 1973, *Astrophys. J.* **186**, 481.

Hohl, F.: 1973, *Astrophys. J.* **184**, 353.

Miller, R. H., Prendergast, K. H., and Quirk, W. J.: 1970, *Astrophys. J.* **161**, 903.

Prendergast, K. H. and Tomer, E.: 1970, *Astron. J.* **75**, 674.

## DISCUSSION

*Miller*: A technical question: recurrence-relations such as you mention are usually unstable, and require great care unless you are careful about starting and the direction of recurrence. How do you get stable results?

*Van Albada*: No direct tests of the behaviour of the recurrence relations have been made. The results of experiments with rotating homogeneous spheres agree well with the 'exact' solutions.

*Ipser*: Have you tried to use this method to construct self-consistent solutions for equilibrium configurations?

*Van Albada*: No.

*Spitzer*: Can you compare the computing time required by your method and the corresponding time required by the method of Dr Gott?

*Van Albada*: For a system with many particles, moving them takes longer than evaluating the forces. Solving Poisson's equation on a grid of 1000 cells takes about 1 s on a CDC Cyber 74-16.