

UNIQUENESS RESULTS FOR A CLASS OF HIGHER – ORDER BOUNDARY VALUE PROBLEMS

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Abstract. The classical maximum principle is utilized to obtain maximum principles for functionals which are defined on solutions of fourth, sixth and eighth-order elliptic equations. The principles derived lead to uniqueness results.

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1. Introduction. It is well known that every subharmonic function in a bounded domain Ω (i.e. $\Delta u \geq 0$ in Ω) satisfies the classical maximum principle

$$\max_{\bar{\Omega}} u = \max_{\partial\Omega} u. \quad (1.1)$$

The subbiharmonic function $u(x) = -x_1^4 - |x|^2$ in the ball $\Omega = \{(x_1, \dots, x_n) \mid |x| < R\}$ (i.e. $\Delta^2 u \leq 0$ in Ω) shows that there are no classical maximum principles for the biharmonic operator $\Delta^2 u$ (and for higher-order elliptic operators at all). Still some results can be proven.

The first proof of a maximum principle for an elliptic equation of higher-order was given by Miranda [4]. The Miranda result was extended to semilinear equations by Schaefer [5] and used to deduce the non-existence of solution to certain boundary value problems. An extension of this last maximum principle to a sixth-order equation may be found in [6]. There are other works dealing with maximum principles for fourth and sixth-order equations and their applications (see [1], [2], [3], [7] (Chapter 10), [9] and the references cited therein). A general maximum principle for an equation with constant coefficients of order $2m$ was established by Tseng and Lin ([8]).

Our purpose here is to deduce some maximum principles for a class of fourth and sixth-order equations which do not appear to be contained in the above mentioned works. A maximum principle for an equation of eighth-order with non-constant coefficients is also presented. From these maximum principles we obtain uniqueness results for a class of higher-order equations in a plane domain Ω . If the domain is on the line, then we briefly indicate some uniqueness results for problems of order $2m$ ($m \geq 5$).

2. Maximum principles and uniqueness results. Let Ω be a bounded domain in the plane with a smooth boundary $\partial\Omega$. We denote partial derivatives $\partial u / \partial x_i$ by $u_{,i}$ and use the summation convention, so that, e.g., the square of the gradient of u becomes

$(\partial u/\partial x_1)^2 + (\partial u/\partial x_2)^2 = |\nabla u|^2 = u_{,i} u_{,i}$. The scalar product of the gradients of two functions, say u and v , may be written as $\nabla u \cdot \nabla v = u_{,i} v_{,i}$.

First, we consider a fourth-order equation

$$\Delta^2 u - a(x)\Delta u + b(x)u = 0 \quad \text{in } \Omega, \tag{2.1}$$

where

$$a > 0, \quad \Delta(1/a) \leq 0 \tag{2.2}$$

and

$$b \geq 0. \tag{2.3}$$

We have the following result.

LEMMA 2.1. *If $u \in C^4(\Omega) \cap C^2(\overline{\Omega})$ is a (classical) solution of (2.1), then the functional F_1 given by*

$$F_1 = \frac{a(x)}{2} u^2 + (|\nabla u|^2 - u\Delta u) \tag{2.4}$$

assumes its maximum value on $\partial\Omega$.

From Lemma 2.1 we now deduce a slight extension of Schaefer’s result (6, Theorem 4).

THEOREM 2.1. *There is at most one classical solution of the boundary value problem*

$$\begin{cases} \Delta^2 u - a(x)\Delta u + b(x)u = f & \text{in } \Omega \\ u = g, \quad \frac{\partial u}{\partial n} = h & \text{on } \partial\Omega, \end{cases} \tag{2.5}$$

where a, b satisfy (2.2) and (2.3) and w is a C^1 function for which $bw' > 0$ in Ω .

Here and throughout the paper $\partial/\partial n$ denotes the outward normal derivative operator.

It is quite obvious how Theorem 1 and Theorem 4 in [5] have to be modified to obtain the above stated results. The details are left to the reader.

We now prove a maximum principle for the sixth-order equation

$$\Delta^3 u - a(x)\Delta^2 u + b(x)\Delta u - c(x)u = 0 \quad \text{in } \Omega \subset \mathbb{R}^2, \tag{2.6}$$

where

$$c > 0, \quad \Delta(1/c) \leq 0. \tag{2.7}$$

LEMMA 2.2. *Let $u \in C^6(\Omega) \cap C^4(\overline{\Omega})$ be a (classical) solution of (2.6) and suppose that (2.2), (2.3) and (2.7) are satisfied. Then the functional F_2 given by*

$$F_2 = \frac{c(x)}{2} u^2 + \frac{a(x)}{2} (\Delta u)^2 + |\nabla(\Delta u)|^2 - \Delta u \Delta^2 u \tag{2.8}$$

assumes its maximum value on $\partial\Omega$.

Proof. The assumption (2.2) implies that

$$\begin{aligned} \Delta(c(x)u^2)/2 &\geq c(x)u\Delta u + \sum_{i=1}^2 \left(\frac{1}{\sqrt{c}} c_{,i} u + \sqrt{c} u_{,i} \right)^2 \\ &\geq c(x)u\Delta u. \end{aligned}$$

Similarly,

$$\Delta(a(x)(\Delta u)^2)/2 \geq a(x)\Delta u\Delta^2 u.$$

Now

$$\begin{aligned} \Delta(|\nabla(\Delta u)|^2) &= 2(\Delta u)_{,ij}(\Delta u)_{,ij} + 2\nabla(\Delta u) \cdot \nabla(\Delta^2 u). \\ \Delta(\Delta u\Delta^2 u) &= (\Delta^2 u)^2 + 2\nabla(\Delta u) \cdot \nabla(\Delta^2 u) + \Delta u\Delta^3 u. \end{aligned}$$

Hence

$$\Delta F_2 \geq 2(\Delta u)_{,ij}(\Delta u)_{,ij} - (\Delta^2 u)^2 + b(x)(\Delta u)^2 \quad \text{in } \Omega.$$

Since the inequality

$$2v_{,ij}v_{,ij} \geq (\Delta v)^2 \tag{2.9}$$

holds in two dimensions, we obtain

$$\Delta F_2 \geq 0 \quad \text{in } \Omega. \quad \square$$

An immediate consequence of Lemma 2.2 is the following uniqueness result.

THEOREM 2.2. *There is at most one classical solution of the boundary value problem*

$$\begin{cases} \Delta^3 u - a(x)\Delta^2 u + b(x)\Delta u - c(x)u = f & \text{in } \Omega \\ u = g, \quad \Delta u = h, \quad \frac{\partial(\Delta u)}{\partial n} = i & \text{on } \partial\Omega, \end{cases} \tag{2.10}$$

where a, b, c satisfy (2.2), (2.3) and (2.7).

Proof. Assume that u_1 and u_2 are solutions of (2.10) and let $v = u_1 - u_2$. The function v satisfies (2.6) and

$$v = \Delta v = \frac{\partial(\Delta v)}{\partial n} = 0 \quad \text{on } \partial\Omega. \tag{2.11}$$

Since $\Delta v = 0$ on $\partial\Omega$ we obtain that $|\nabla(\Delta v)|^2 = \left(\frac{\partial(\Delta v)}{\partial n}\right)^2$ on $\partial\Omega$. Hence, by Lemma 2.2 and (2.11),

$$F_2 \leq \max_{\partial\Omega} F_2 = 0 \quad \text{in } \Omega. \tag{2.12}$$

Consequently,

$$-\Delta v\Delta^2 v \leq 0 \quad \text{in } \Omega. \tag{2.13}$$

Integrating (2.13) over Ω and using Green’s identity we conclude that

$$|\nabla(\Delta v)| = 0 \quad \text{in } \Omega,$$

and by continuity $\Delta v \equiv 0$ in Ω . Since $v = 0$ on Ω we are led to the conclusion that $v \equiv 0$ in Ω . □

Next we deal with classical solutions (i.e. $C^8(\Omega) \cap C^6(\overline{\Omega})$) of

$$\Delta^4 u - a(x)\Delta^3 u + b(x)\Delta^2 u - c(x)\Delta u + du = 0 \quad \text{in } \Omega \subset \mathbb{R}^2, \tag{2.14}$$

where

$$d > 0. \tag{2.15}$$

A uniqueness result can be inferred from the following maximum principle

LEMMA 2.3. *Let u be a classical solution of (2.14). Assume that (2.2), (2.3), (2.7) and (2.15) are satisfied. Then the functional*

$$F_3 = \frac{c(x)}{2}(\Delta u)^2 + \frac{a(x)}{2}(\Delta^2 u)^2 + d(|\nabla u|^2 - u\Delta u) + |\nabla(\Delta^2 u)|^2 - \Delta^2 u\Delta^3 u \tag{2.16}$$

assumes its maximum value on $\partial\Omega$.

Proof. Using (2.2) and (2.7), we have

$$\begin{aligned} \Delta(c(x)(\Delta u)^2)/2 &\geq c(x)\Delta u\Delta^2 u, \\ \Delta(a(x)(\Delta^2 u)^2)/2 &\geq a(x)\Delta^2 u\Delta^3 u. \end{aligned}$$

By inequality (2.9) we get

$$\begin{aligned} d\Delta(|\nabla u|^2 - u\Delta u) &= d(2u_{,ij}u_{,ij} - (\Delta u)^2 - u\Delta^2 u) \geq -du\Delta^2 u, \\ \Delta(|\nabla(\Delta^2 u)|^2 - \Delta^2 u\Delta^3 u) &= 2(\Delta^2 u)_{,ij}(\Delta^2 u)_{,ij} - (\Delta^3 u)^2 - \Delta^2 u\Delta^4 u \geq -\Delta^2 u\Delta^4 u. \end{aligned}$$

Hence from equation (2.14) we obtain

$$\Delta F_3 \geq 0 \quad \text{in } \Omega. \tag{2.17}$$

THEOREM 2.3. *There is at most one classical solution of the boundary value problem*

$$\begin{cases} \Delta^4 u - a\Delta^3 u + b(x)\Delta^2 u - c\Delta u + du = f & \text{in } \Omega \\ u = g, \quad \Delta u = h, \quad \Delta^2 u = i, \quad \frac{\partial(\Delta^2 u)}{\partial n} = j & \text{on } \partial\Omega, \end{cases} \tag{2.17}$$

where a satisfies (2.2), b satisfies (2.3), c satisfies (2.7), $d > 0$, and the curvature k of $\partial\Omega$ is positive.

Proof. To establish this result, we suppose that u_1 and u_2 are two solutions of (2.17). Defining $v = u_1 - u_2$, we see that v satisfies (2.14) and

$$v = \Delta v = \Delta^2 v = \frac{\partial(\Delta^2 v)}{\partial n} = 0 \quad \text{on } \partial\Omega. \tag{2.18}$$

Since $v = \Delta^2 v = 0$ on $\partial\Omega$, we obtain

$$\begin{aligned} \frac{\partial F_3}{\partial n} &= c\Delta v \frac{\partial(\Delta v)}{\partial n} + a\Delta^2 v \frac{\partial(\Delta^2 v)}{\partial n} + d \left(2 \frac{\partial v}{\partial n} \frac{\partial^2 v}{\partial n^2} - \frac{\partial v}{\partial n} \Delta v - v \frac{\partial(\Delta v)}{\partial n} \right) \\ &\quad + 2 \frac{\partial(\Delta^2 v)}{\partial n} \frac{\partial^2(\Delta^2 v)}{\partial n^2} - \frac{\partial(\Delta^2 v)}{\partial n} \Delta^3 v - \Delta^2 v \frac{\partial(\Delta^3 v)}{\partial n} \quad \text{on } \partial\Omega. \end{aligned}$$

By (2.18) we have

$$\frac{\partial F_3}{\partial n} = 2d \frac{\partial v}{\partial n} \frac{\partial^2 v}{\partial n^2} \quad \text{on } \partial\Omega. \tag{2.19}$$

By introducing normal coordinates in the neighbourhood of the boundary, we can write

$$\Delta v = \frac{\partial^2 v}{\partial n^2} + \frac{\partial^2 v}{\partial s^2} + k \frac{\partial v}{\partial n}, \tag{2.20}$$

where $\frac{\partial v}{\partial s}$ denotes the tangential derivative of v . Since $v = \Delta v = 0$ on $\partial\Omega$, relation (2.20) becomes

$$\frac{\partial^2 v}{\partial n^2} = -k \frac{\partial v}{\partial n},$$

which gives

$$\frac{\partial F_3}{\partial n} = -2dk \left(\frac{\partial v}{\partial n} \right)^2 \leq 0 \quad \text{on } \partial\Omega.$$

This contradicts Hopf’s lemma at a point $P \in \partial\Omega$, where F_3 ($F_3 \neq \text{constant}$) assumes its maximum value. Hence F_3 is constant in $\bar{\Omega}$. Thus $\frac{\partial F_3}{\partial n} = 0$ on $\partial\Omega$ and consequently $\frac{\partial v}{\partial n} = 0$ on $\partial\Omega$. By the boundary conditions it follows that $F_3 = 0$ in $\bar{\Omega}$. Using Green’s identity we arrive at $v = 0$ in $\bar{\Omega}$. □

Finally we shift our attention from the two dimensional to the one dimensional case and mention the following result (Ω denotes an open interval (α, β)).

THEOREM 2.4. *There can be at most one classical solution of the problem*

$$\begin{cases} u^{(2m)} - du^{(6)} + c(x)u^{(4)} - b(x)u'' + a(x)u = f & \text{in } \Omega \\ u = g_1, u'' = g_2, u''' = g_3, \dots, u^{(m)} = g_m & \text{on } \partial\Omega, \end{cases} \tag{2.21}$$

where $m \geq 6$ is even, $d \geq 0$ and $b \geq 0$, $a, c > 0$, $(1/a)''$, $(1/c)'' \leq 0$ in Ω .

The result follows since the function

$$\begin{aligned} F_4 &= u''u^{(2m-2)} - 2u'''u^{(2m-3)} + 3u^{(4)}u^{(2m-4)} - \dots + (m-3)u^{(m-2)}u^{(m+2)} \\ &\quad - (m-3)u^{(m-1)}u^{(m+1)}/2 - ((m-3)/2 + 1)[(u^{(m)})^2 - u^{(m-1)}u^{(m+1)}] \\ &\quad + [(u''')^2 - du''u^{(4)}] + c(x)(u'')^2/2 + a(x)u^2/2 \end{aligned}$$

assumes its maximum value on $\partial\Omega$, where u is a solution of

$$u^{(2m)} - du^{(6)} + c(x)u^{(4)} - b(x)u'' + a(x)u = 0 \quad \text{in } \Omega.$$

Similarly, we can treat the problem

$$\begin{cases} u^{(2m)} + du^{(6)} - c(x)u^{(4)} + b(x)u'' - a(x)u = f & \text{in } \Omega \\ u = g_1, u' = g_2, u'' = g_3, \dots, u^{(m)} = g_m & \text{on } \partial\Omega, \end{cases} \quad (2.22)$$

where $m \geq 5$ is odd.

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