



Relations for quadratic Hodge integrals via stable maps

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Abstract. Following Faber–Pandharipande, we use the virtual localization formula for the moduli space of stable maps to \mathbb{P}^1 to compute relations between Hodge integrals. We prove that certain generating series of these integrals are polynomials.

1 Introduction

Let $\overline{\mathcal{M}}_{g,n}$ be the moduli space of n -pointed genus g stable curves. It is a proper smooth Deligne Mumford (DM) stack of dimension $3g - 3 + n$. We denote by $\pi : \overline{\mathcal{C}}_{g,n} \rightarrow \overline{\mathcal{M}}_{g,n}$ the universal curve and by $\sigma_i : \overline{\mathcal{M}}_{g,n} \rightarrow \overline{\mathcal{C}}_{g,n}$ the sections associated with the marking i for all $1 \leq i \leq n$. We denote by $\omega_{\overline{\mathcal{C}}_{g,n}/\overline{\mathcal{M}}_{g,n}}$ the relative dualizing sheaf of π . We will consider the following classes in $A^*(\overline{\mathcal{M}}_{g,n})$:

- For all $0 \leq i \leq g$, λ_i stands for the i th Chern class of the Hodge bundle, i.e., the vector bundle $\mathbb{E} = \pi_* \omega_{\overline{\mathcal{C}}_{g,n}/\overline{\mathcal{M}}_{g,n}}$. For all $\alpha \in \mathbb{C}$, we denote $\Lambda_g(\alpha) = \sum_{j=0}^g \alpha^{g-j} \lambda_j$, and $\Lambda_g^\vee(\alpha) = (-1)^g \Lambda_g(-\alpha)$.¹
- For all $1 \leq i \leq n$, we denote ψ_i the Chern class of the cotangent line at the i th marking $\mathcal{L}_i = \sigma_i^*(\omega_{\overline{\mathcal{C}}_{g,n}/\overline{\mathcal{M}}_{g,n}})$.

A *Hodge integral* is an intersection number of the form:

$$\int_{\overline{\mathcal{M}}_{g,n}} \psi_1^{k_1} \dots \psi_n^{k_n} \Lambda_g(t_1) \dots \Lambda_g(t_m),$$

where k_1, \dots, k_n are nonnegative integers and t_1, \dots, t_m are complex numbers. If $m = 1, 2$, or 3 , then the above integral is called a linear, double, or triple Hodge integrals, respectively. Relations between linear Hodge integrals were proved in [FP00a] using the Gromov–Witten theory of \mathbb{P}^1 and the localization formula of [GP99]. This approach was also used in [FP00b] and [TZ03] to prove certain properties of triple Hodge integrals. Linear and triple Hodge integrals naturally appeared in the GW-theory of Calabi–Yau 3-folds, thus explaining a more abundant literature on the topic. However, double Hodge integrals have appeared recently in the quantization of Witten–Kontsevich generating series (see [Blo20]), in the theory of spin Hurwitz

Received by the editors June 12, 2023; revised January 8, 2024; accepted January 9, 2024.

Published online on Cambridge Core January 17, 2024.

AMS subject classification: 14H10.

Keywords: Hodge integrals, moduli of curves, intersection theory, localization formula.

¹Here, we use the convention of [FP00a] for $\Lambda_g^\vee(\alpha)$ and $\Lambda_g(\alpha)$.



numbers (see [GKL21]), and in the GW theory of blow-ups of smooth surfaces (see [GKLS22]).

In the present note, we consider the following power series in $\mathbb{C}[\alpha][[t]]$ defined using double Hodge integrals:

$$P_a(\alpha, t) = \sum_{g \geq 0} t^g \left(\int_{\overline{\mathcal{M}}_{g,n+1}} \frac{\Lambda_g^\vee(1)\Lambda_g^\vee(\alpha)}{1 - \psi_0} \prod_{i=1}^n (2a_i + 1)!! (-4\psi_i)^{a_i} \right) \exp\left(\frac{t}{24}\right),$$

where $a = (a_1, \dots, a_n)$ is a vector of nonnegative integers. If $n = 1$, we use the convention: $\int_{\overline{\mathcal{M}}_{0,2}} \psi_1^a \frac{\Lambda_g^\vee(1)\Lambda_g^\vee(\alpha)}{1 - \psi_2} = (-1)^a$.

Theorem 1.1 $P_a(\alpha, t)$ is a monic polynomial in $\mathbb{C}[\alpha][t]$ of degree $|a|$ in t .

Here, we provide the first values of $P_a(-\alpha - 1, t)$. In the list below, we omit the variables $-\alpha - 1$ and t in the notation:

$$P_{()} = 1.$$

$$P_{(1)} = t + 12.$$

$$P_{(2)} = t^2 - 10\alpha t + 240.$$

$$P_{(1,1)} = t^2 - 12t.$$

$$P_{(3)} = t^3 + (-77/3\alpha - 28)t^2 + 280t + 6720.$$

$$P_{(2,1)} = t^3 + (-10\alpha - 48)t^2 + (240\alpha + 240)t.$$

$$P_{(1,1,1)} = t^3 - 72t^2 + 432t.$$

$$P_{(4)} = t^4 + (-43\alpha - 72)t^3 + (126\alpha^2 + 756\alpha + 840)t^2 + 10080t + 241920.$$

$$P_{(3,1)} = t^4 + (-77/3\alpha - 100)t^3 + (1232\alpha + 1624)t^2.$$

$$P_{(2,2)} = t^4 + (20\alpha + 100)t^3 + (-100\alpha^2 - 1360\alpha - 1680)t^2.$$

$$P_{(2,1,1)} = t^4 + (-10\alpha - 132)t^3 + (840\alpha + 3120)t^2 + (-8640\alpha - 8640)t.$$

$$P_{(1,1,1,1)} = t^4 - 168t^3 + 5616t^2 - 20736t.$$

Considering these first values, we conjecture that P_a is a polynomial of total degree $|a|$ in both variables t and α .

2 Preliminaries

We denote by $\overline{\mathcal{M}}_{g,n}(\mathbb{P}^1, 1)$, the moduli space of stable maps of degree 1 to \mathbb{P}^1 . It is a proper DM stack of virtual dimension $2g + n$. Here, we can define in an analogous way the Hodge bundle \mathbb{E} , the cotangent line bundles \mathcal{L}_i and we denote again λ_i and ψ_i the respective Chern classes. We also have the forgetful and evaluation maps

$$\pi: \overline{\mathcal{M}}_{g,n+1}(\mathbb{P}^1, 1) \rightarrow \overline{\mathcal{M}}_{g,n}(\mathbb{P}^1, 1), \text{ and } ev_i: \overline{\mathcal{M}}_{g,n+1}(\mathbb{P}^1, 1) \rightarrow \mathbb{P}^1.$$

Throughout this note, the enumeration of markings starts from 0. Furthermore, π is the morphism that forgets the marking p_0 and ev_i is the evaluation of a stable map to the i th marked point. The vector bundle $T := R^1\pi_*(ev_0^*\mathcal{O}_{\mathbb{P}^1}(-1))$ is of rank g and we denote by y its top Chern class. We will denote:

$$\langle \prod_{i=0}^{n-1} \tau_{a_i}(\omega) | y \rangle_{g,1}^{\mathbb{P}^1} := \int_{[\overline{\mathcal{M}}_{g,n}(\mathbb{P}^1,1)]^{vir}} \prod_{i=0}^{n-1} \psi_i^{a_i} ev_i^*(\omega) y,$$

where ω denotes the class of a point in \mathbb{P}^1 .

Theorem 2.1 (Localization Formula [GP99, FP00a]) *Let $g \in \mathbb{Z}_{\geq 0}$, and let $a \in \mathbb{Z}_{\geq 0}^n$ such that $|a| \leq g$. Then, for all complex numbers α , and $t \in \mathbb{C}^*$, we have*

$$\begin{aligned} \langle \prod_{i=1}^n \tau_{a_i}(\omega) | y \rangle_{g,1}^{\mathbb{P}^1} &= \sum_{g_1+g_2=g} \int_{\overline{\mathcal{M}}_{g_1,n+1}} t^n \prod_{i=1}^n \psi_i^{a_i} \frac{\Lambda_{g_1}^\vee(t) \Lambda_{g_1}^\vee(\alpha t)}{t(t-\psi_0)} \\ &\quad \times \int_{\overline{\mathcal{M}}_{g_2,1}} \frac{\Lambda_{g_2}^\vee(-t) \Lambda_{g_2}^\vee((\alpha+1)t)}{-t(-t-\psi_0)}. \end{aligned}$$

Here, we use the convention $\int_{\overline{\mathcal{M}}_{0,1}} \psi_0^a = 1$.

Proposition 2.2 (Proposition 4.1 of [TZ03]) *For all complex numbers α , we have*

$$F(\alpha, t) = 1 + \sum_{g>0} t^{2g} \int_{\overline{\mathcal{M}}_{g,1}} \frac{\Lambda_g^\vee(1) \Lambda_g^\vee(\alpha)}{1-\psi_0} = \exp\left(-\frac{t^2}{24}\right).$$

Besides, we have the String and Dilaton equation for Hodge integrals.

Proposition 2.3 *Let $g, n \in \mathbb{Z}_{\geq 0}$ such that $2g - 2 + n > 0$.*

(i) *[Dilaton equation for Hodge integrals] Let $(a_1, \dots, a_n) \in \mathbb{Z}_{\geq 0}^n$ and assume that there exist i_0 such that $a_{i_0} = 1$. Then*

$$\int_{\overline{\mathcal{M}}_{g,n+1}} \frac{\psi_{i_0} \prod_{i \neq i_0} \psi_i^{a_i} \prod_{j=1}^g \lambda_k^{b_k}}{1-\psi_0} = (2g-2+n) \int_{\overline{\mathcal{M}}_{g,n}} \frac{\prod_{i=1}^{n-1} \psi_i^{a_i} \prod_{j=1}^g \lambda_k^{b_k}}{1-\psi_0}.$$

(ii) *[String equation for Hodge integrals] Let $(a_1, \dots, a_n) \in \mathbb{Z}_{\geq 0}^n$ and assume that there exist i_0 such that $a_{i_0} = 0$. Then we have*

$$\begin{aligned} \int_{\overline{\mathcal{M}}_{g,n+1}} \frac{\prod_{i=1}^n \psi_i^{a_i} \prod_{j=1}^g \lambda_k^{b_k}}{1-\psi_0} &= \int_{\overline{\mathcal{M}}_{g,n}} \frac{\prod_{i=1}^{n-1} \psi_i^{a_i} \prod_{j=1}^g \lambda_k^{b_k}}{1-\psi_0} \\ &\quad + \sum_{j=1}^n \int_{\overline{\mathcal{M}}_{g,n}} \frac{\psi_j^{a_j-1} \prod_{i \neq j} \psi_i^{a_i} \prod_{k=1}^g \lambda_k^{b_k}}{1-\psi_0}. \end{aligned}$$

3 The calculation

Note that the GW-invariant $\langle \prod_{i=1}^n \tau_{a_i}(\omega) | y \rangle_{g,1}^{\mathbb{P}^1}$ is 0 unless $|a| = g$ for dimensional reasons. Indeed, $\dim_{\mathbb{C}}[\overline{\mathcal{M}}_{g,n}(\mathbb{P}^1,1)]^{vir} = 2g+n$ and the cycle we are integrating is in codimension $g+|a|+n$. Using the above localization formula, and Lemma 2.1 of

[TZ03] the intersection number $\langle \prod_{i=1}^n \tau_{a_i}(\omega) | y \rangle_{g,1}^{\mathbb{P}^1}$ is expressed as

$$\begin{aligned} & \sum_{g_1+g_2=g} \int_{\overline{\mathcal{M}}_{g_1,n+1}} t^n \prod_{i=1}^n \psi_i^{a_i} \frac{\Lambda_{g_1}^\vee(t) \Lambda_{g_1}^\vee(\alpha t)}{t(t-\psi_0)} \cdot \int_{\overline{\mathcal{M}}_{g_2,1}} \frac{\Lambda_{g_2}^\vee(-t) \Lambda_{g_2}^\vee((\alpha+1)t)}{-t(-t-\psi_0)} \\ &= \sum_{g_1+g_2=g} t^{|a|-g_1} (-t)^{-g_2} \int_{\overline{\mathcal{M}}_{g_1,n+1}} \prod_{i=1}^n \psi_i^{a_i} \frac{\Lambda_{g_1}^\vee(1) \Lambda_{g_1}^\vee(\alpha)}{1-\psi_0} \times \int_{\overline{\mathcal{M}}_{g_2,1}} \frac{\Lambda_{g_2}^\vee(1) \Lambda_{g_2}^\vee(-(\alpha+1))}{1-\psi_0} \\ &= t^{|a|-g} \sum_{g_1+g_2=g} \int_{\overline{\mathcal{M}}_{g_1,n+1}} \prod_{i=1}^n \psi_i^{a_i} \frac{\Lambda_{g_1}^\vee(1) \Lambda_{g_1}^\vee(\alpha)}{1-\psi_0} \cdot \int_{\overline{\mathcal{M}}_{g_2,1}} \psi_0^{3g_2-2}. \end{aligned}$$

In the last equation, we used Proposition 2.2 in order to replace $\int_{\overline{\mathcal{M}}_{g_2,1}} \frac{\Lambda_{g_2}^\vee(1) \Lambda_{g_2}^\vee(-(\alpha+1))}{1-\psi_0}$ with $(-1)^{g_2} \int_{\overline{\mathcal{M}}_{g_2,1}} \psi_0^{3g_2-2}$.

We define

$$A_{g,a}(\alpha) = \sum_{g_1+g_2=g} \int_{\overline{\mathcal{M}}_{g_1,n+1}} \prod_{i=1}^n \psi_i^{a_i} \frac{\Lambda_{g_1}^\vee(1) \Lambda_{g_1}^\vee(\alpha)}{1-\psi_0} \cdot \int_{\overline{\mathcal{M}}_{g_2,1}} \psi_0^{3g_2-2}.$$

Then, we have

$$A_{g,a}(\alpha) = \begin{cases} 0, & |a| < g, \\ \langle \prod_{i=1}^n \tau_{a_i}(\omega) | y \rangle_{g,1}^{\mathbb{P}^1}, & |a| = g. \end{cases}$$

By the definition of $\Lambda_g^\vee(t)$, we see that $\Lambda_g^\vee(1) \Lambda_g^\vee(-(\alpha+1))$ is a polynomial in α of degree g , which actually determines the degree of $A_g(\alpha)$.

We now present a proof for the main result.

Proof (of Theorem 1.1) We begin by stating the well-known fact

$$1 + \sum_{g>0} t^g \int_{\overline{\mathcal{M}}_{g,1}} \psi_0^{3g-2} = \exp\left(\frac{t}{24}\right)$$

proven in Section 3.1 of [FP00a]. Now, we consider the product of $\exp\left(\frac{t}{24}\right)$ and

$$\sum_{g \geq 0} t^g \left(\int_{\overline{\mathcal{M}}_{g,n+1}} \frac{\Lambda_g^\vee(1) \Lambda_g^\vee(\alpha)}{1-\psi_0} \prod_{i=1}^n (2a_i + 1)!! (-4\psi_i)^{a_i} \right)$$

to obtain a new power series whose coefficients in degree g are given by

$$\sum_{g_1+g_2=g} \int_{\overline{\mathcal{M}}_{g_1,n+1}} \prod_{i=1}^n (2a_i + 1)!! (-4)^{a_i} \prod_{i=1}^n \psi_i^{a_i} \frac{\Lambda_{g_1}^\vee(1) \Lambda_{g_1}^\vee(\alpha)}{1-\psi_0} \cdot \int_{\overline{\mathcal{M}}_{g_2,1}} \psi_0^{3g_2-2}.$$

This is exactly $A_{g,a}(\alpha) \cdot \prod_{i=1}^n (2a_i + 1)!! (-4)^{a_i}$. Hence, we can rewrite the power series $P_a(\alpha, t)$ in the form

$$P_a(\alpha, t) = \prod_{i=1}^n (2a_i + 1)!! (-4)^{a_i} \sum_{g \geq 0} t^g A_{g,a}(\alpha).$$

As it is computed in the start of Section 3, we have that the numbers $A_{g,a}(\alpha)$ vanish when $g > |a|$. Hence, we get that all coefficients of the power series $P_a(\alpha, t)$

vanish when $g > |a|$, i.e. $P_a(\alpha, t)$ is a polynomial of degree $|a|$. Furthermore, the top coefficient of $P_a(\alpha, t)$, i.e., the coefficient of $t^{|a|}$ is given by

$$\left\langle \prod_{i=1}^n (-4)^{a_i} (2a_i + 1)!! \tau_{a_i}(\omega) | \gamma \right\rangle_{|a|, 1}^{\mathbb{P}^1}.$$

This value is computed in [KL1] and is actually equal to 1. In particular, the number $\prod_{i=1}^n (-4)^{a_i} (2a_i + 1)!!$ is here to make the polynomial monic. ■

We now prove several other properties of the polynomials P_a .

Proposition 3.1 *The constant term c_0 of $P_a(\alpha, t)$ is nonzero if and only if $n = 1$, where then $c_0 = (-1)^a \prod_{i=1}^n (-4)^{a_i} (2a_i + 1)!!$ or if $n > 1$ and $\sum_{i=1}^n a_i \leq n - 2$ where then*

$$c_0 = \prod_{i=1}^n (-4)^{a_i} (2a_i + 1)!! \frac{(n - 2)!}{a_1! \dots (n - 2 - \sum a_i)!}.$$

Proof We only compute the integrals appearing in the constant term of this polynomial since then we only have to multiply with $\prod_{i=1}^n (2a_i + 1)!! (-4)^{a_i}$. The integral in the constant term of $P_a(\alpha, t)$ is given by $\int_{\overline{\mathcal{M}}_{0, n+1}} \frac{\prod_{i=1}^n \psi_i^{a_i}}{1 - \psi_0}$. When $n = 1$, using the convention $\int_{\overline{\mathcal{M}}_{0, 2}} \frac{\psi_1^a}{1 - \psi_0} = (-1)^a$, we get that

$$c_0 = (-1)^a \prod_{i=1}^n (-4)^{a_i} (2a_i + 1)!!.$$

When $n > 1$, if $\sum_{i=1}^n a_i > n - 2$, then c_0 is zero for dimensional reasons. Otherwise, we have

$$\int_{\overline{\mathcal{M}}_{0, n+1}} \frac{\prod_{i=1}^n \psi_i^{a_i}}{1 - \psi_0} = \int_{\overline{\mathcal{M}}_{0, n+1}} \psi_0^{n-2-\sum a_i} \prod_{i=1}^n \psi_i^{a_i} = \frac{(n - 2)!}{a_1! \dots (n - 2 - \sum a_i)!}. \quad \blacksquare$$

Proposition 3.2 *Let $n \geq 3$. Then we have the following rules:*

(i) [String equation]

$$P_{(a_1, \dots, a_{n-1}, 0)}(\alpha, t) = P_{(a_1, \dots, a_{n-1})}(\alpha, t) - \sum_{i=1}^n (8a_i + 4) P_{(a_1, \dots, a_i-1, \dots, a_{n-1})}(\alpha, t).$$

(ii) [Dilaton equation]

$$P_{(a_1, \dots, a_{n-1}, 1)}(\alpha, t) = (t - 12n + 24) P_{(a_1, \dots, a_{n-1})}(\alpha, t) - 24t P'_{(a_1, \dots, a_{n-1})}(\alpha, t).$$

Proof We define the power series

$$\tilde{P}_a(\alpha, t) = \sum_{g \geq 0} t^g \left(\int_{\overline{\mathcal{M}}_{g, n+1}} \prod_{i=1}^n \psi_i^{a_i} \frac{\Lambda_g^\vee(1) \Lambda_g^\vee(\alpha)}{1 - \psi_0} \right).$$

Note that the following equation holds:

$$P_a(\alpha, t) = \prod_{i=1}^n (2a_i + 1)!! (-4)^{a_i} \tilde{P}_a(\alpha, t) \exp\left(\frac{t}{24}\right).$$

We can rewrite the coefficients of $\tilde{P}_a(\alpha, t)$ as

$$\sum_{k=0}^g \sum_{j=0}^g (-1)^{g+k} (a+1)^{g-j} \int_{\overline{\mathcal{M}}_{g,n+1}} \frac{\prod_{i=1}^n \psi_i^{a_i} \lambda_k \lambda_j}{1 - \psi_0}.$$

(i) Applying the String equation for Hodge integrals, we obtain the following formula:

$$\tilde{P}_{(a_1, \dots, a_{n-1}, 0)}(\alpha, t) = \tilde{P}_{(a_1, \dots, a_{n-1})}(\alpha, t) + \sum_{i=1}^n \tilde{P}_{(a_1, \dots, a_i-1, \dots, a_{n-1})}(\alpha, t).$$

Hence, multiplying with $\prod_{i=1}^{n-1} (2a_i + 1)!! (-4)^{a_i} \exp\left(\frac{t}{24}\right)$, we obtain the desired result after a straightforward calculation.

(ii) Applying Dilaton equation for Hodge integrals, we obtain the following formula:

$$\begin{aligned} \tilde{P}_{(a_1, \dots, a_{n-1}, 1)}(\alpha, t) &= 2 \sum_{g \geq 0} g t^g \int_{\overline{\mathcal{M}}_{g,n-1}} \prod_{i=1}^{n-1} \psi_i^{a_i} \frac{\Lambda_g^\vee(1) \Lambda_g^\vee(\alpha)}{1 - \psi_0} \\ &\quad + (n-2) \tilde{P}_{(a_1, \dots, a_{n-1})}(\alpha, t). \end{aligned}$$

Note that the first term of the sum is equal to $2t \tilde{P}'_{(a_1, \dots, a_{n-1})}(\alpha, t)$. Now, multiplying both sides of the equation above with

$$\prod_{i=1}^{n-1} (2a_i + 1) (-4)^{a_i} \exp\left(\frac{t}{24}\right),$$

we have

$$\begin{aligned} \frac{-1}{12} P_{(a_1, \dots, a_{n-1}, 1)}(\alpha, t) &= (n-2) P_{(a_1, \dots, a_{n-1})}(\alpha, t) \\ &\quad + 2t \left(\prod_{i=1}^{n-1} (-4)^{a_i} (2a_i + 1)!! \right) \tilde{P}'_{(a_1, \dots, a_{n-1})}(\alpha, t) e^{t/24} \\ &= (n-2) P_{(a_1, \dots, a_{n-1})}(\alpha, t) \\ &\quad + 2t (P'_{(a_1, \dots, a_{n-1})}(\alpha, t) - \frac{1}{24} P_{(a_1, \dots, a_{n-1})}(\alpha, t)). \end{aligned}$$

Finally, clearing the denominators, we obtain the desired result. ■

We recall Mumford’s relation $\Lambda_g^\vee(1) \cdot \Lambda_g^\vee(-1) = 1$ (see [Mum83]). In particular, $P_a(-1, t)$ is defined by integrals of ψ -classes.

Corollary 3.3 For any vector $a \in \mathbb{Z}_{\geq 0}^n$, the power series

$$P_a(-1, t) = \prod_{i=1}^n (2a_i + 1)!! (-4)^{a_i} \exp\left(\frac{t}{24}\right) \cdot \sum_{g \geq 0} (-t)^g \int_{\overline{\mathcal{M}}_{g,n+1}} \frac{\prod_{i=1}^n \psi_i^{a_i}}{1 - \psi_0}$$

is a polynomial of degree $|a|$.

In this case, the polynomiality as well as a closed expression were proved in [LX11].

Acknowledgments I am very grateful to my PhD advisor Adrien Sauvaget for introducing me to this problem and for his guidance and comments throughout the whole writing of this article.

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